# Differentiaal Vergelijkingen In de Aardwetenschappen PDEs - chapt 13

C. Thieulot (c.thieulot@uu.nl)

January 2015

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

► Laplace's equation

$$\nabla^2 u = 0$$

<□ > < @ > < E > < E > E のQ @

► Laplace's equation

$$\nabla^2 u = 0$$

Poisson's equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} = f(\boldsymbol{x}, \boldsymbol{y}, \boldsymbol{z})$$

Laplace's equation

$$\nabla^2 u = 0$$

Poisson's equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} = f(\boldsymbol{x}, \boldsymbol{y}, \boldsymbol{z})$$

diffusion equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} = \frac{1}{\alpha^2} \frac{\partial \boldsymbol{u}}{\partial t}$$

▲□▶ ▲圖▶ ▲圖▶ ▲圖▶ = ● ● ●

Laplace's equation

$$\nabla^2 u = 0$$

Poisson's equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} = f(\boldsymbol{x}, \boldsymbol{y}, \boldsymbol{z})$$

diffusion equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} = \frac{1}{\alpha^2} \frac{\partial \boldsymbol{u}}{\partial t}$$

wave equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} = \frac{1}{\boldsymbol{v}^2} \frac{\partial^2 \boldsymbol{u}}{\partial t^2}$$

Laplace's equation

$$\nabla^2 u = 0$$

Poisson's equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} = f(\boldsymbol{x},\boldsymbol{y},\boldsymbol{z})$$

diffusion equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} = \frac{1}{\alpha^2} \frac{\partial \boldsymbol{u}}{\partial t}$$

wave equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} = \frac{1}{\boldsymbol{v}^2} \frac{\partial^2 \boldsymbol{u}}{\partial t^2}$$

Helmholtz equation

 $\boldsymbol{\nabla}^2 \boldsymbol{u} + \boldsymbol{k}^2 \boldsymbol{F} = \boldsymbol{0}$ 

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

Laplace's equation

$$\nabla^2 u = 0$$

Poisson's equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} = f(\boldsymbol{x}, \boldsymbol{y}, \boldsymbol{z})$$

diffusion equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} = \frac{1}{\alpha^2} \frac{\partial \boldsymbol{u}}{\partial t}$$

wave equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} = \frac{1}{\boldsymbol{v}^2} \frac{\partial^2 \boldsymbol{u}}{\partial t^2}$$

Helmholtz equation

$$\boldsymbol{\nabla}^2 \boldsymbol{u} + \boldsymbol{k}^2 \boldsymbol{F} = \boldsymbol{0}$$

Schrödinger equation

$$-\frac{\hbar^2}{2m}\nabla^2\Psi + V\Psi = i\hbar\frac{\partial\Psi}{\partial t}$$

$$\nabla^2 u = \Delta u = 0$$

no dependence on time, just on the spatial variables.

$$\nabla^2 u = \Delta u = 0$$

- no dependence on time, just on the spatial variables.
- Laplace Equation describes steady state situations such as :
  - steady state temperature distributions
  - steady state stress distributions
  - steady state potential distributions (it is also called the potential equation

steady state flows, for example in a cylinder, around a corner

In cartesian coordinates :

▶ 1D :

$$\frac{\partial^2 u}{\partial x^2} = 0$$

In cartesian coordinates :

▶ 1D :

$$\frac{\partial^2 u}{\partial x^2} = 0$$

► 2D :

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$

In cartesian coordinates :

 $\frac{\partial^2 u}{\partial x^2} = 0$ 

► 2D :  $\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$ 

▶ 1D :

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} = 0$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

Example 1 : (Maxwell's equations)

The electric field is related to the charge density by the divergence relationship

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

with

- **E** : electric field
- $\rho$  : charge density
- $\epsilon_0$  : vacuum permittivity (=8.854187817620... × 10<sup>-12</sup> Farads per metre)

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

Example 1 : (Maxwell's equations)

The electric field is related to the charge density by the divergence relationship

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

with

- ► E : electric field
- $\rho$  : charge density
- $\epsilon_0$  : vacuum permittivity (=8.854187817620... × 10<sup>-12</sup> Farads per metre)

The electric field is related to the electric potential by a gradient relationship

 $\mathbf{E} = -\nabla V$ 

Example 1 : (Maxwell's equations)

The electric field is related to the charge density by the divergence relationship

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

with

► E : electric field

•  $\rho$  : charge density

•  $\epsilon_0$ : vacuum permittivity (=8.854187817620... × 10<sup>-12</sup> Farads per metre) The electric field is related to the electric potential by a gradient relationship

$$\mathbf{E} = - \mathbf{\nabla} V$$

In a charge-free region of space ( $\rho=$  0), the potential is therefore related to the charge density by Laplace equation

$$\nabla^2 V = 0$$

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

Example 2 : The heat transport equation writes :

$$\rho c_{p} \left( \frac{\partial T}{\partial t} + \underbrace{\mathbf{v} \cdot \nabla T}_{adv.} \right) = \underbrace{\nabla \cdot (k \nabla T)}_{diff.} + \underbrace{H}_{prod.}$$

Example 2 : The heat transport equation writes :

$$\rho c_{p} \left( \frac{\partial T}{\partial t} + \underbrace{\mathbf{v} \cdot \nabla T}_{adv.} \right) = \underbrace{\nabla \cdot (k \nabla T)}_{diff.} + \underbrace{H}_{prod.}$$

Assume :

• steady state 
$$(\partial/\partial t \rightarrow 0)$$

Example 2 : The heat transport equation writes :

$$\rho c_{p} \left( \frac{\partial T}{\partial t} + \underbrace{\mathbf{v} \cdot \nabla T}_{adv.} \right) = \underbrace{\nabla \cdot (k \nabla T)}_{diff.} + \underbrace{H}_{prod.}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

Assume :

- steady state  $(\partial/\partial t \rightarrow 0)$
- no advection  $(\mathbf{v} = \mathbf{0})$

Example 2 :

The heat transport equation writes :

$$\rho c_{\rho} \left( \frac{\partial T}{\partial t} + \underbrace{\mathbf{v} \cdot \nabla T}_{adv.} \right) = \underbrace{\nabla \cdot (k \nabla T)}_{diff.} + \underbrace{H}_{prod.}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

Assume :

- steady state  $(\partial/\partial t \rightarrow 0)$
- no advection  $(\mathbf{v} = \mathbf{0})$
- no radiogenic heat production (H = 0)

Example 2 :

The heat transport equation writes :

$$\rho c_{\rho} \left( \frac{\partial T}{\partial t} + \underbrace{\mathbf{v} \cdot \nabla T}_{adv.} \right) = \underbrace{\nabla \cdot (k \nabla T)}_{diff.} + \underbrace{H}_{prod.}$$

Assume :

- steady state  $(\partial/\partial t \rightarrow 0)$
- no advection  $(\mathbf{v} = \mathbf{0})$
- no radiogenic heat production (H = 0)

Then the temperature field must verify :

$$\boldsymbol{\nabla}\cdot(k\boldsymbol{\nabla}T)=0$$

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

Example 2 :

The heat transport equation writes :

$$\rho c_{p} \left( \frac{\partial T}{\partial t} + \underbrace{\mathbf{v} \cdot \nabla T}_{adv.} \right) = \underbrace{\nabla \cdot (k \nabla T)}_{diff.} + \underbrace{H}_{prod.}$$

Assume :

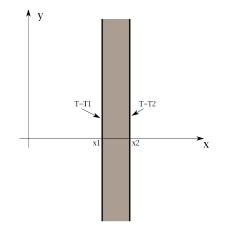
- steady state  $(\partial/\partial t \rightarrow 0)$
- no advection  $(\mathbf{v} = \mathbf{0})$
- no radiogenic heat production (H = 0)

Then the temperature field must verify :

$$\boldsymbol{\nabla}\cdot(k\boldsymbol{\nabla}T)=0$$

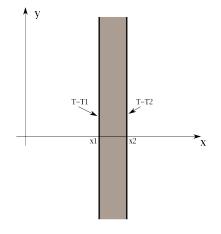
If k = constant, then it is equivalent to solve

$$\nabla^2 T = 0$$



▲□▶ ▲圖▶ ★ 国▶ ★ 国▶ - 国 - のへで

plate is infinite in y and z direction :  $\frac{\partial}{\partial y} \to 0$  ,  $\to \frac{\partial}{\partial z} \to 0$ 



▲□▶ ▲圖▶ ▲臣▶ ▲臣▶ ―臣 … のへで

plate is infinite in y and z direction :  $\frac{\partial}{\partial y} \to 0$ ,  $\to \frac{\partial}{\partial z} \to 0$  $\Rightarrow T$  is a function of x only.

We have to solve :

$$\frac{d^2 T}{dx^2} = 0$$

with boundary conditions :

$$T(x_1) = T1 \qquad T(x_2) = T2$$

(ロ)、(型)、(E)、(E)、 E) の(の)

We have to solve :

$$\frac{d^2 T}{dx^2} = 0$$

with boundary conditions :

$$T(x_1) = T1 \qquad T(x_2) = T2$$

We integrate once :

$$\frac{dT}{dx} = a$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 のへぐ

We have to solve :

$$\frac{d^2 T}{dx^2} = 0$$

with boundary conditions :

$$T(x_1) = T1 \qquad T(x_2) = T2$$

We integrate once :

$$\frac{dT}{dx} = a$$

and another time :

$$T(x) = ax + b$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

We have to solve :

$$\frac{d^2 T}{dx^2} = 0$$

with boundary conditions :

$$T(x_1) = T1 \qquad T(x_2) = T2$$

We integrate once :

$$\frac{dT}{dx} = a$$

and another time :

$$T(x) = ax + b$$

We use the b.c. to determine a and b:

$$T(x_1) = T_1 = ax_1 + b$$
  
 $T(x_2) = T_2 = ax_2 + b$ 

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

We have to solve :

$$\frac{d^2 T}{dx^2} = 0$$

with boundary conditions :

$$T(x_1) = T1 \qquad T(x_2) = T2$$

We integrate once :

$$\frac{dT}{dx} = a$$

and another time :

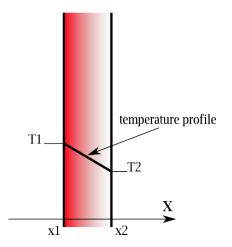
$$T(x) = ax + b$$

We use the b.c. to determine a and b:

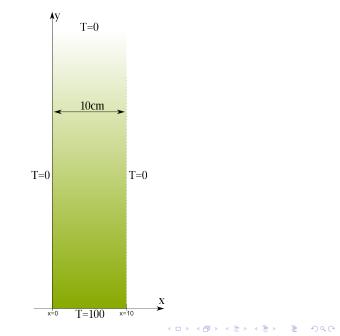
$$T(x_1) = T_1 = ax_1 + b$$
  
 $T(x_2) = T_2 = ax_2 + b$ 

leading to  $a = \frac{T_2 - T_1}{x_2 - x_1}$  and  $b = T_1 - \frac{T_2 - T_1}{x_2 - x_1}x_1$  and finally

$$T(x) = \frac{T_2 - T_1}{x_2 - x_1}(x - x_1) + T_1$$



◆□> ◆□> ◆豆> ◆豆> ・豆 ・ のへで



The temperature satisfies the 2D Laplace equation inside the plate :

$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = 0 \tag{1}$$

(ロ)、(型)、(E)、(E)、 E) の(の)

The temperature satisfies the 2D Laplace equation inside the plate :

$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = 0 \tag{1}$$

We could try to solve the equation by using a tentative solution of the form :

$$T(x,y) = \theta(x)\Phi(y)$$
(2)

Be do not *know* the solution is of this form.

The temperature satisfies the 2D Laplace equation inside the plate :

$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = 0 \tag{1}$$

We could try to solve the equation by using a tentative solution of the form :

$$T(x,y) = \theta(x)\Phi(y)$$
(2)

Be do not *know* the solution is of this form.

We substitute (2) into (1) and obtain :

$$\Phi \frac{\partial^2 \theta}{\partial x^2} + \theta \frac{\partial^2 \Phi}{\partial y^2} = 0$$

The temperature satisfies the 2D Laplace equation inside the plate :

$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = 0 \tag{1}$$

We could try to solve the equation by using a tentative solution of the form :

$$T(x,y) = \theta(x)\Phi(y)$$
(2)

Be do not *know* the solution is of this form.

We substitute (2) into (1) and obtain :

$$\Phi \frac{\partial^2 \theta}{\partial x^2} + \theta \frac{\partial^2 \Phi}{\partial y^2} = 0$$

Dividing by  $\theta \Phi$  gives :

$$\frac{1}{\theta}\frac{\partial^2\theta}{\partial x^2} + \frac{1}{\Phi}\frac{\partial^2\Phi}{\partial y^2} = 0$$

$$\frac{1}{\theta}\frac{\partial^2\theta}{\partial x^2} + \frac{1}{\Phi}\frac{\partial^2\Phi}{\partial y^2} = 0$$

<□ > < @ > < E > < E > E のQ @

$$\frac{1}{\theta}\frac{\partial^2\theta}{\partial x^2} + \frac{1}{\Phi}\frac{\partial^2\Phi}{\partial y^2} = 0$$

Separation of variables : we say that each term is a constant because the first term is a function of x only and the second a function of y only. We then write

$$\frac{1}{\theta}\frac{\partial^2\theta}{\partial x^2} = -\frac{1}{\Phi}\frac{\partial^2\Phi}{\partial y^2} = -k^2$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

where k is called the separation constant.

$$\frac{1}{\theta}\frac{\partial^2\theta}{\partial x^2} + \frac{1}{\Phi}\frac{\partial^2\Phi}{\partial y^2} = 0$$

Separation of variables : we say that each term is a constant because the first term is a function of x only and the second a function of y only. We then write

$$\frac{1}{\theta}\frac{\partial^2\theta}{\partial x^2} = -\frac{1}{\Phi}\frac{\partial^2\Phi}{\partial y^2} = -k^2$$

where k is called the separation constant. This leads to

$$\frac{\partial^2 \theta}{\partial x^2} + k^2 \theta = 0$$
$$\frac{\partial^2 \Phi}{\partial y^2} - k^2 \Phi = 0$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

$$\frac{1}{\theta}\frac{\partial^2\theta}{\partial x^2} + \frac{1}{\Phi}\frac{\partial^2\Phi}{\partial y^2} = 0$$

Separation of variables : we say that each term is a constant because the first term is a function of x only and the second a function of y only. We then write

$$\frac{1}{\theta}\frac{\partial^2\theta}{\partial x^2} = -\frac{1}{\Phi}\frac{\partial^2\Phi}{\partial y^2} = -k^2$$

where k is called the separation constant. This leads to

$$\frac{\partial^2 \theta}{\partial x^2} + k^2 \theta = 0$$
$$\frac{\partial^2 \Phi}{\partial y^2} - k^2 \Phi = 0$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへで

• The solution to the first one is  $\theta(x) = \sin kx$  or  $\theta(x) = \cos kx$ 

$$\frac{1}{\theta}\frac{\partial^2\theta}{\partial x^2} + \frac{1}{\Phi}\frac{\partial^2\Phi}{\partial y^2} = 0$$

Separation of variables : we say that each term is a constant because the first term is a function of x only and the second a function of y only. We then write

$$\frac{1}{\theta}\frac{\partial^2\theta}{\partial x^2} = -\frac{1}{\Phi}\frac{\partial^2\Phi}{\partial y^2} = -k^2$$

where k is called the separation constant. This leads to

$$\frac{\partial^2 \theta}{\partial x^2} + k^2 \theta = 0$$
$$\frac{\partial^2 \Phi}{\partial y^2} - k^2 \Phi = 0$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ □ □ のへで

- The solution to the first one is  $\theta(x) = \sin kx$  or  $\theta(x) = \cos kx$
- The solution to the second one is  $\Phi(x) = e^{kx}$  or  $\Phi(x) = e^{-kx}$

The general solution writes :

$$T(x,y) = \theta(x)\Phi(y) = \left\{ \begin{array}{c} \sin kx \\ \cos kx \end{array} \right\} \left\{ \begin{array}{c} e^{ky} \\ e^{-ky} \end{array} \right\}$$

The general solution writes :

$$T(x,y) = \theta(x)\Phi(y) = \begin{cases} \sin kx \\ \cos kx \end{cases} \begin{cases} e^{ky} \\ e^{-ky} \end{cases}$$

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶

We can now use the b.c. to find the solution to the Laplace equation.

The general solution writes :

$$T(x,y) = \theta(x)\Phi(y) = \left\{\begin{array}{c} \sin kx \\ \cos kx \end{array}\right\} \left\{\begin{array}{c} e^{ky} \\ e^{-ky} \end{array}\right\}$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへで

We can now use the b.c. to find the solution to the Laplace equation.

• Since  $T \to 0$  when  $y \to \infty$  then  $e^{ky}$  unacceptable.

The general solution writes :

$$T(x,y) = \theta(x)\Phi(y) = \left\{ egin{array}{c} \sin kx \ \cos kx \end{array} 
ight\} \left\{ egin{array}{c} e^{ky} \ e^{-ky} \end{array} 
ight\}$$

< ロ > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < 回 > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □

We can now use the b.c. to find the solution to the Laplace equation.

- Since  $T \to 0$  when  $y \to \infty$  then  $e^{ky}$  unacceptable.
- Since T = 0 when x = 0 then  $\cos kx$  unacceptable.

The general solution writes :

$$T(x,y) = \theta(x)\Phi(y) = \begin{cases} \sin kx \\ \cos kx \end{cases} \begin{cases} e^{ky} \\ e^{-ky} \end{cases}$$

We can now use the b.c. to find the solution to the Laplace equation.

- Since  $T \to 0$  when  $y \to \infty$  then  $e^{ky}$  unacceptable.
- Since T = 0 when x = 0 then  $\cos kx$  unacceptable.

so

$$T(x,y)=\sin(kx)\ e^{-ky}$$

The general solution writes :

$$T(x,y) = \theta(x)\Phi(y) = \left\{ egin{array}{c} \sin kx \ \cos kx \end{array} 
ight\} \left\{ egin{array}{c} e^{ky} \ e^{-ky} \end{array} 
ight\}$$

We can now use the b.c. to find the solution to the Laplace equation.

- Since  $T \to 0$  when  $y \to \infty$  then  $e^{ky}$  unacceptable.
- Since T = 0 when x = 0 then  $\cos kx$  unacceptable.

so

$$T(x,y) = \sin(kx) e^{-ky}$$

We finally use T = 0 at x = 10 which leads to  $10k = n\pi$ , i.e. :

$$T(x,y) = \sin(\frac{n\pi x}{10}) e^{-n\pi y/10}$$

Problem : the solution does not satisfy T(x, 0) = 100 !

Problem : the solution does not satisfy T(x,0) = 100 !

PANIC A linear combination of solutions is still a solution !

Problem : the solution does not satisfy T(x, 0) = 100 !

PANC A linear combination of solutions is still a solution !

Let's find such a combination which satisfies the b.c. at y = 0:

$$T(x,y) = \sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi x}{10}\right) e^{-n\pi y/10}$$

Problem : the solution does not satisfy T(x,0) = 100 !

PANC A linear combination of solutions is still a solution !

Let's find such a combination which satisfies the b.c. at y = 0:

$$T(x,y) = \sum_{n=1}^{\infty} b_n \sin(\frac{n\pi x}{10}) e^{-n\pi y/10}$$

We impose then T(x,0) = 100:

$$100 = \sum_{n=1}^{\infty} b_n \sin(\frac{n\pi x}{10})$$

$$100 = \sum_{n=1}^{\infty} b_n \sin(\frac{n\pi x}{10})$$

This is the Fourier sine series of f(x) = 100 with l = 10 (chapter 7.9 of Boas).

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□▶

$$100 = \sum_{n=1}^{\infty} b_n \sin(\frac{n\pi x}{10})$$

This is the Fourier sine series of f(x) = 100 with l = 10 (chapter 7.9 of Boas). The coefficient  $b_n$  is then given by

$$b_n = \frac{2}{l} \int_0^l f(x) \sin \frac{n\pi x}{l} dx = \frac{2}{10} \int_0^l 100 \sin \frac{n\pi x}{10} dx = \begin{cases} 400/n\pi & \text{odd n} \\ 0 & \text{even n} \end{cases}$$

◆□▶ ◆□▶ ◆臣▶ ◆臣▶ 臣 の�?

$$100 = \sum_{n=1}^{\infty} b_n \sin(\frac{n\pi x}{10})$$

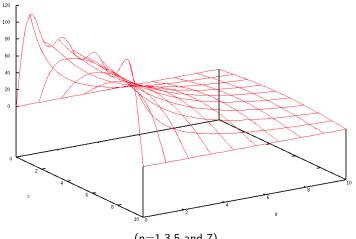
This is the Fourier sine series of f(x) = 100 with l = 10 (chapter 7.9 of Boas). The coefficient  $b_n$  is then given by

$$b_n = \frac{2}{l} \int_0^l f(x) \sin \frac{n\pi x}{l} dx = \frac{2}{10} \int_0^l 100 \sin \frac{n\pi x}{10} dx = \begin{cases} 400/n\pi & \text{odd n} \\ 0 & \text{even n} \end{cases}$$

Finally (!) :

$$T(x,y) = \frac{400}{\pi} \left( e^{-\pi y/10} \sin(\frac{\pi x}{10}) + \frac{1}{3} \sin(\frac{3\pi x}{10}) e^{-3\pi y/10} + \dots \right)$$





(n=1,3,5 and 7)

・ロト ・ 日 ・ ・ 日 ・ ・ 日 ・ æ

## Numerical solution (1)

The heat transport equation writes :

$$\rho c_{\rho} \left( \frac{\partial T}{\partial t} + \underbrace{\mathbf{v} \cdot \nabla T}_{adv.} \right) = \underbrace{\nabla \cdot (k \nabla T)}_{diff.} + \underbrace{H}_{prod.}$$

▲□▶ ▲圖▶ ▲圖▶ ▲圖▶ = ● ● ●

## Numerical solution (1)

The heat transport equation writes :

$$\rho c_{p} \left( \frac{\partial T}{\partial t} + \underbrace{\mathbf{v} \cdot \nabla T}_{adv.} \right) = \underbrace{\nabla \cdot (k \nabla T)}_{diff.} + \underbrace{H}_{prod.}$$

▲□▶ ▲圖▶ ▲圖▶ ▲圖▶ = ● ● ●

Assume :

• no advection  $(\mathbf{v} = \mathbf{0})$ 

#### Numerical solution (1)

The heat transport equation writes :

$$\rho c_{p} \left( \frac{\partial T}{\partial t} + \underbrace{\mathbf{v} \cdot \nabla T}_{adv.} \right) = \underbrace{\nabla \cdot (k \nabla T)}_{diff.} + \underbrace{H}_{prod.}$$

Assume :

- no advection  $(\mathbf{v} = \mathbf{0})$
- no radiogenic heat production (H = 0)

Then the temperature field must verify :

$$\rho c_p \frac{\partial T}{\partial t} = \underbrace{\nabla \cdot (k \nabla T)}_{\text{diff.}}$$

And at steady state ( $\partial_t = 0$ )

$$\boldsymbol{\nabla} \cdot (\boldsymbol{\nabla} T) = 0$$

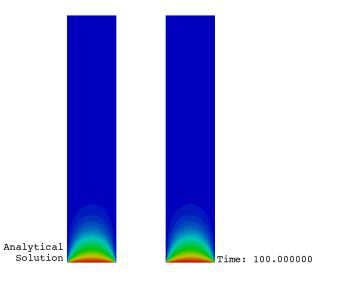
◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

# Numerical solution (2)

▲□▶ ▲圖▶ ▲圖▶ ▲圖▶ = ● ● ●

film

# Numerical solution (3)



◆□ → ◆□ → ◆三 → ◆三 → ◆□ →